Monod-Wyman-Changeux Analysis of Ligand-Gated Ion Channel Mutants

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Supporting Information

ABSTRACT: We present a framework for computing the gating properties of ligand-gated ion channel mutants using the Monod-Wyman-Changeux (MWC) model of allostery. We derive simple analytic formulas for key functional properties such as the leakiness, dynamic range, half-maximal effective concentration ([EC50]), and effective Hill coefficient, and explore the full spectrum of phenotypes that are accessible through mutations. Specifically, we consider mutations in the channel pore of nicotinic acetylcholine receptor (nAChR) and the ligand binding domain of a cyclic nucleotide-gated (CNG) ion channel, demonstrating how each mutation can be characterized as only affecting a subset of the biophysical parameters. In addition, we show how the unifying perspective offered by the MWC model allows us, perhaps surprisingly, to collapse the plethora of dose-response data from different classes of ion channels into a universal family of curves.

1. INTRODUCTION

Ion channels are signaling proteins responsible for a huge variety of physiological functions ranging from responding to membrane voltage, tension, and temperature to serving as the primary players in the signal transduction we experience as vision.1 Broadly speaking, these channels are classified on the basis of the driving forces that gate them. In this work, we explore one such classification for ligand-gated ion channel mutants based on the Monod-Wyman-Changeux (MWC) model of allostery. In particular, we focus on mutants in two of the arguably best-studied ligand-gated ion channels: the nicotinic acetylcholine receptor (nAChR) and the cyclic nucleotide-gated (CNG) ion channel shown schematically in Figure 1.2,3

The MWC model has long been used in the contexts of both nAChR and CNG ion channels.4−6 Although careful analysis of these systems has revealed that some details of ligand-gated ion channel dynamics are not captured by this model (e.g., the existence and interaction of multiple conducting states7,8), the MWC framework nevertheless captures many critical features of a channel’s response and offers one of the simplest settings to explore its underlying mechanisms. For example, knowledge of both the nAChR and CNG systems’ molecular architecture and our ability to measure their complex kinetics has only recently become sufficiently advanced to tie the effects of mutations to key biophysical parameters. Purohit and Auerbach used combinations of mutations to infer the nAChR gating energy, finding that unliganded nAChR channels open up for a remarkably brief 80 μs every 15 min.9 Statistical mechanics has been used to show how changes to the energy difference between conformations in allosteric proteins translate to different functional behavior (i.e., how it modifies the leakiness, dynamic range, [EC50], and the effective Hill coefficient),10,11 and we extend this work to find simple analytic approximations that are valid within the context of ion channels. Using this methodology, we systematically explore the full range of behaviors that may be induced by different types of mutations. This analysis enables us to quantify the inherent trade-offs between key properties of ion channel dose-response curves and potentially paves the way for future biophysical models of evolutionary fitness in which the genotype (i.e., amino acid sequence) of allosteric molecules is directly connected to phenotype (i.e., properties of a channel’s response).

To this end, we consider two distinct classes of mutants which tune different sets of MWC parameters, either the ion channel gating energy or the ligand-channel dissociation constants. Previous work by Auerbach et al. demonstrated that these two sets of physical parameters can be independently tuned within the nAChR ion channel; pore mutations only alter the channel gating energy whereas mutations within the ligand
binding domain only affect the ligand-channel dissociation constants. Utilizing this parameter independence, we determine the full spectrum of nAChR phenotypes given an arbitrary set of channel pore mutations and show why a linear increase in the channel gating energy leads to a logarithmic shift in the nAChR dose-response curve. Next, we consider recent data from CNGA2 ion channels with mutations in their ligand binding pocket. We hypothesize that modifying the ligand binding domain should not alter the channel gating energy and demonstrate how the entire class of CNGA2 mutants can be...
**MODEL**

Electrophysiological techniques can measure currents across a single cell’s membrane. The current flowing through a ligand-gated ion channel is proportional to the average probability $p_{open}(c)$ that the channel will be open at a ligand concentration $c$. For an ion channel with $m$ identical ligand binding sites (see Figure 2), this probability is given by the MWC model as

$$p_{open}(c) = \frac{\left(1 + \frac{c}{K_C}\right)^m}{\left(1 + \frac{c}{K_O}\right)^m + e^{-\beta c} \left(1 + \frac{c}{K_C}\right)^m}$$

(1)

where $K_O$ and $K_C$ represent the dissociation constants between the ligand and the open and closed ion channel, respectively, $c$ denotes the concentration of the ligand, $e$ (called the gating energy) denotes the free energy difference between the closed and open conformations of the ion channel in the absence of ligand, and $\beta = \frac{1}{k_B T}$ where $k_B$ is Boltzmann’s constant and $T$ is the temperature of the system. Wild-type ion channels are typically closed in the absence of ligand ($e < 0$) and open when bound to ligand ($K_O < K_C$). Figure 2 shows the possible conformations of the nAChR ($m = 2$) and CNGA2 ($m = 4$) ion channels together with their Boltzmann weights. $p_{open}(c)$ is given by the sum of the open state weights divided by the sum of all weights. Note that the MWC parameters $K_O$, $K_C$, and $e$ may be expressed as ratios of the experimentally measured rates of ligand binding and unbinding as well as the transition rates between the open and closed channel conformations (see Supporting Information section A.1).

Current measurements are often reported as normalized current, implying that the current has been stretched vertically to run from 0 to 1, as given by

$$\text{normalized current} = \frac{p_{open}(c) - p_{open}^{\min}}{p_{open}^{\max} - p_{open}^{\min}}$$

(2)

$p_{open}(c)$ increases monotonically as a function of ligand concentration $c$, with a minimum value in the absence of ligand given by

$$p_{open}^{\min} = p_{open}(0) = \frac{1}{1 + e^{-\beta c}}$$

(3)

and a maximum value in the presence of saturating levels of ligand given as

$$p_{open}^{\max} = \lim_{c \to \infty} p_{open}(c) = \frac{1}{1 + e^{-\beta c} \left(\frac{K_C}{K_O}\right)^m}$$

(4)

Using the above two limits, we can investigate four important characteristics of ion channels. First, we examine the leakiness of an ion channel, or the fraction of time a channel is open in the absence of ligand, namely

$$\text{leakiness} = p_{open}^{\min}$$

(5)

Next we determine the dynamic range, or the difference between the probability of the maximally open and maximally closed states of the ion channel, given by

$$\text{dynamic range} = p_{open}^{\max} - p_{open}^{\min}$$

(6)

Ion channels that minimize leakiness only open upon ligand binding, and ion channels that maximize dynamic range have greater contrast between their open and closed states. Just like $p_{open}(c)$, leakiness and dynamic range lie within the interval $[0, 1]$.

Two other important characteristics are measured from the normalized current. The half-maximal effective concentration $[EC_{50}]$ denotes the concentration of ligand at which the normalized current of the ion channel equals 1/2, namely

$$p_{open}([EC_{50}]) = \frac{p_{open}^{\min} + p_{open}^{\max}}{2}$$

(7)

The effective Hill coefficient $h$ equals twice the log-log slope of the normalized current evaluated at $c = [EC_{50}]$

$$h = \frac{d}{d \log c} \log \left(\frac{p_{open}(c) - p_{open}^{\min}}{p_{open}^{\max} - p_{open}^{\min}}\right)_{c=[EC_{50}]}$$

(8)

which reduces to the standard Hill coefficient for the Hill function. The $[EC_{50}]$ determines how the normalized current shifts left and right, while the effective Hill coefficient corresponds to the slope at $[EC_{50}]$. Together, these two properties determine the approximate window of ligand concentrations for which the normalized current transitions from 0 to 1.

In the limit $1 \ll e^{-\beta c} \ll \left(\frac{K_C}{K_O}\right)^m$, which we show below is relevant for both the nAChR and CNGA2 ion channels, the various functional properties of the channel described above can be approximated to leading order as (see Supporting Information section B):

$$\text{leakiness} \approx e^{\beta c}$$

(9)

$$\text{dynamic range} \approx 1$$

(10)

$$[EC_{50}] \approx e^{-\beta c/m} K_O$$

(11)

$$h \approx m$$

(12)
**RESULTS**

nAChR Mutants Can Be Categorized Using Free Energy. Muscle-type nAChR is a heteropentamer with subunit stoichiometry α₂βγδ, containing two ligand binding sites for acetylcholine at the interface of the α-δ and α-γ subunits.

The five homologous subunits have M2 transmembrane domains which move symmetrically during nAChR gating to either occlude or open the ion channel. By introducing a serine in place of the leucine at a key residue (L251S) within the M2 domain present within each subunit, the corresponding subunit is able to more easily transition from the closed to open configuration, shifting the dose-response curve to the left (see Figure 3A). For example, wild-type nAChR is maximally activated by the agonist acetylcholine (ACh), while a mutant ion channel with one L251S mutation is more sensitive and only requires 10M of acetylcholine to saturate its dose-response curve.

Focusing on how the L251S mutation causes a shift in the dose-response curve, we use the normalized current data for several of these mutants; from right to left, the curves represent n = 0 (wild-type) to n = 4 (an ion channel with four of its five subunits mutated). One interesting trend in the data is that each additional mutation shifts the normalized current to the left by approximately one decade in concentration, highlighting the trade-offs between the physical parameters of the system (in the case of nAChR, between the entropy associated with the ligand concentration and the gating free energy). When the Bohr parameters of two ion channels are equal, both channels will elicit the same gating energy, and the probability of the receptor being open is given by the following equation:

\[ P_{open}(c) = \frac{1}{1 + e^{-\beta F(c)}} \]  

where \( F(c) = -k_B T \log \left( \frac{1 + c K_C}{1 + c K_0} \right) \)  

The Bohr parameter quantifies the trade-offs between the physical parameters of the system (in the case of nAChR, between the entropy associated with the ligand concentration and the gating free energy). The Bohr parameter maps the full complexity of a generic ion channel response into a single combination of the relevant physical parameters of the system.

Full Spectrum of nAChR Gating Energy Mutants. We next consider the entire range of nAChR phenotypes achievable by modifying the gating free energy \( \epsilon \) of the wild-type ion channel. For instance, any combination of nAChR pore mutations would be expected to not affect the ligand dissociation constants and thus yield an ion channel within this class (see Supporting Information section A.4 for one such example). For concreteness, we focus on how the \( \epsilon \) parameter tunes key features of the dose-response curves, namely, the leftness, dynamic range, \( [EC_{50}] \), and effective Hill coefficient \( h \).
Figure 4. Theoretical prediction and experimental measurements for mutant nAChR ion channel characteristics. The open squares mark the $\beta \epsilon$ values of the five dose-response curves from Figure 3A. (A) The leakiness given by eq 5 increases exponentially with each mutation. (B) The dynamic range from eq 6 is nearly uniform for all mutants. (C) The $[EC_{50}]$ decreases exponentially with each mutation. (D) The effective Hill coefficient $h$ is predicted to remain approximately constant. $[EC_{50}]$ and $h$ offer a direct comparison between the best-fit model predictions (□) and the experimental measurements (●) from Figure 3A. While the $[EC_{50}]$ matches well between theory and experiment, the effective Hill coefficient $h$ is significantly noisier.

(see eqs 5–12), although we note that other important phenotypic properties such as the intrinsic noise and capacity have also been characterized for the MWC model. Figure 4 shows these four characteristics, with the open squares representing the properties of the five best-fit dose-response curves from Figure 3A.

Figure 4A implies that all of the mutants considered here have negligible leakiness; the probability that the wild-type channel ($\beta \epsilon(0) = -23.7$) will be open is less than $10^{-10}$. Experimental measurements have shown that such spontaneous openings occur extremely infrequently in nAChR, although direct measurement is difficult for such rare events. Other mutational analysis has predicted gating energies around $\beta \epsilon(0) \approx -14$ (corresponding to a leakiness of $10^{-5}$), but we note that such a large wild-type gating energy prohibits the five mutants in Figure 3 from being fit as a single class of mutants with the same $K_O$ and $K_C$ values (see Supporting Information section C.2). If this large wild-type gating energy is correct, it may imply that the L251S mutation also affects the $K_O$ and $K_C$ parameters, though the absence of error bars on the original data make it hard to quantitatively assess the underlying origins of these discrepancies.

Figure 4B asserts that all of the mutant ion channels should have full dynamic range except for the wild-type channel, which has a dynamic range of 0.91. In comparison, the measured dynamic range of wild-type nAChR is 0.95, close to our predicted value. Note that only when the dynamic range approaches unity does the normalized current become identical to $p_{\text{open}}$; for lower values, information about the leakiness and dynamic range is lost by only measuring normalized currents.

We compare the $[EC_{50}]$ (Figure 4C) and effective Hill coefficient $h$ (Figure 4D) with the nAChR data by interpolating the measurements (see Supporting Information section C.3) in order to precisely determine the midpoint and slope of the response. The $[EC_{50}]$ predictions faithfully match the data over 4 orders of magnitude. Because each additional mutation lowers the $[EC_{50}]$ by approximately one decade, the analytic form $\epsilon + k_B T$ implies that $\epsilon$ increases by roughly $5k_B T$ per mutation, enabling the ion channel to open more easily. In addition to the L251S mutation considered here, another mutation (L251T) has also been found to shift $[EC_{50}]$ by a constant logarithmic amount (see Supporting Information section A.4). We also note that many biological systems logarithmically tune their responses by altering the energy difference between two allosteric states, as seen through processes such as phosphorylation and calmodulin binding. This may give rise to an interesting interplay between physiological time scales where such processes occur and evolutionary time scales where traits such as the $[EC_{50}]$ may be accessed via mutations like those considered here.

Lastly, the Hill coefficients of the entire class of mutants all lie between 1.5 and 2.0 except for the $n = 3$ mutant whose dose-response curve in Figure 3A is seen to be flatter than the MWC prediction. We also note that if the L251S mutation moderately perturbs the $K_O$ and $K_C$ values, it would permit fits that more finely attune to the specific shape of each mutant’s data set. That said, the dose-response curve for the $n = 3$ mutant could

\[ p_{\text{open}} = 10^{-12} \]

\[ p_{\text{open}} = 10^{-10} \]
easily be shifted by small changes in the measured values, and hence, without recourse to error bars, it is difficult to make definitive statements about the value adopted for this mutant.

Note that the simplified expressions eqs 9–12 for the leakiness, dynamic range, $[EC_{50}]$, and effective Hill coefficient apply when $1 \ll e^{-\beta \epsilon} \ll \left(\frac{K_O}{K_C}\right)^m$, which given the values of $K_C$ and $K_C$ for the nAChR mutant class translates to $-22 \lesssim \beta \epsilon \lesssim -5$. The $n = 1, 2$, and 3 mutants all fall within this range, and hence, each subsequent mutation exponentially increases their leakiness and exponentially decreases their $[EC_{50}]$, while their dynamic range and effective Hill coefficient remain indifferent to the L251S mutation. The $\beta \epsilon$ parameters of the $n = 0$ and $n = 4$ mutants lie at the edge of the region of validity, so higher order approximations can be used to more precisely fit their functional characteristics (see Supporting Information section B).

**Hetero-Oligomeric CNGA2 Mutants Can Be Categorized Using an Expanded MWC Model.** The nAChR mutant class discussed above had two equivalent ligand binding sites, and only the gating free energy $\beta \epsilon$ varied for the mutants we considered. In this section, we use beautiful data for the olfactory CNGA2 ion channel to explore the unique phenotypes that emerge from a hetero-oligomeric ion channel whose subunits have different ligand binding strengths.

The wild-type CNGA2 ion channel is made up of four identical subunits, each with one binding site for the cyclic nucleotide ligands cAMP or cGMP.26 Within the MWC model, the probability that this channel is open is given by eq 1 with $m = 4$ ligand binding sites (see Figure 2B). Wongsamitkul et al. constructed a mutated subunit with lower affinity for ligand and Figure 5. States and weights for mutant CNGA2 ion channels. CNGA2 mutants with $m = 4$ subunits were constructed using $n$ mutated (light red) and $m - n$ wild-type subunits (purple). The affinity between the wild-type subunits to ligand in the open and closed states ($K_O$ and $K_C$) is stronger than the affinity of the mutated subunits ($K_O^*$ and $K_C^*$). The weights shown account for all possible ligand configurations, with the inset explicitly showing all of the closed states for the wild-type ($n = 0$) ion channel from Figure 2B. The probability that a receptor with $n$ mutated subunits is open is given by its corresponding open state weight divided by the sum of open and closed weights in that same row.

**Figure 6.** Normalized currents for CNGA2 ion channels with a varying number $n$ of mutant subunits. (A) Dose-response curves for CNGA2 mutants composed of $4 - n$ wild-type subunits and $n$ mutated subunits with weaker affinity for the ligand cGMP.13 Once the free energy $\epsilon$ and the ligand dissociation constants of the wild-type subunits ($K_O$ and $K_C$) and mutated subunits ($K_O^*$ and $K_C^*$) are fixed, each mutant is completely characterized by the number of mutated subunits $n$ in eq 15. Theoretical best-fit curves are shown using the parameters $K_O = 1.2 \times 10^{-6}$ M, $K_C = 20 \times 10^{-6}$ M, $K_O^* = 500 \times 10^{-6}$ M, $K_C^* = 140 \times 10^{-3}$ M, and $\beta \epsilon = -3.4$. (B) Data from all five mutants collapses onto a single master curve when plotted as a function of the Bohr parameter given by eq 13. See Supporting Information section C for details on the fitting.
formed tetrameric CNGA2 channels from different combinations of mutated and wild-type subunits (see Figure S). Since the mutation specifically targeted the ligand binding sites, these mutant subunits were postulated to have new ligand dissociation constants but the same free energy difference $\beta$.

We can extend the MWC model to compute the probability $p_{\text{open}}(c)$ that these CNGA2 constructs will be open. The states and weights of an ion channel with $n$ mutated subunits (with ligand affinities $K^\beta_i$ and $K^\epsilon_i$) and $m-n$ wild-type subunits (with ligand affinities $K_O$ and $K_C$) are shown in Figure S, and its probability to be open is given by

$$p_{\text{open}}(c) = \frac{(1 + \frac{c}{K_O})^m(1 + \frac{c}{K_C})^n}{(1 + \frac{c}{K_O})^m(1 + \frac{c}{K_C})^n + e^{-\beta c}(1 + \frac{c}{K^\beta_0})^m(1 + \frac{c}{K^\epsilon_0})^n}$$

(15)

Measurements have confirmed that the dose-response curves of the mutant CNGA2 channels only depend on the total number of mutated subunits $n$ and not on the positions of those subunits (for example, both $n = 2$ with adjacent mutant subunits and $n = 2$ with mutant subunits on opposite corners have identical dose-response curves).

Figure 6A shows the normalized current of all five CNGA2 constructs fit to a single set of $K_O$, $K_C$, $K^\beta_0$, $K^\epsilon_0$, and $c$ parameters. Since the mutated subunits have weaker affinity to ligand (leading to the larger dissociation constants $K^\beta_0 > K_O$ and $K^\epsilon_0 > K_C$), the [EC_{50}] shifts to the right as $n$ increases. As in the case of nAChR, we can collapse the data from this family of mutants onto a single master curve using the Bohr parameter $F_{\text{CNGA2}}(c)$ from eqs 13 and 15, as shown in Figure 6B.

Although we analyze the CNGA2 ion channels in equilibrium, we can glimpse the dynamic nature of the system by computing the probability of each channel conformation. Figure 7A shows the 10 possible states of the wild-type ($n = 0$) channel, the five open states $O_j$ and the five closed states $C_j$ with $0 \leq j \leq 4$ ligands bound. Figure 7B shows how the probabilities of these states are all significantly shifted to the right in the fully mutated ($n = 4$) channel since the mutation diminishes the channel-ligand affinity. The individual state probabilities help determine which of the intermediary states can be ignored when modeling. One extreme simplification that is often made is to consider the Hill limit, where all of the states are ignored save for the closed, unbound ion channel ($C_n$) and the open, fully bound ion channel ($O_n$). The drawbacks of such an approximation are 2-fold: (1) at intermediate ligand concentrations ($c \in [10^{-7}, 10^{-5}]$ M for $n = 0$ and $c \in [10^{-4}, 10^{-2}]$ M for $n = 4$) the ion channel spends at least 10% of its time in the remaining states which results in fundamentally different dynamics than what is predicted by the Hill response, and (2) even in the limits such as $c = 0$ and $c \to \infty$ where the $C_n$ and $O_n$ states dominate the system, the Hill limit ignores the leakiness and dynamic range of the ion channel (requiring them to exactly equal 0 and 1, respectively), thereby glossing over these important properties of the system.

Characterizing CNGA2 Mutants Based on Subunit Composition. We now turn to the leakiness, dynamic range, [EC_{50}], and effective Hill coefficient $h$ of a CNGA2 ion channel with $n$ mutated and $m-n$ wild-type subunits. Detailed derivations for the following results are provided in Supporting Information section B:

$$\text{leakiness} = \frac{1}{1 + e^{-\beta c}}$$

(16)

$$\text{dynamic range} = \frac{1}{1 + e^{-\beta c}(K^\beta_0/K_C)^m/(K^\epsilon_0/K_C)^n} - \frac{1}{1 + e^{-\beta c}}$$

(17)

$$\text{[EC}_{50}\text{]} \approx \begin{cases} e^{-\beta c/m\times K_O} & n = 0 \\ e^{-\beta c/m\times K_O K^*_C/(K^\beta_0 + K_O)} & n = m/2 \\ e^{-\beta c/m\times K^*_C} & n = m \end{cases}$$

(18)

$$h \approx \begin{cases} m/2 & n = m/2 \\ m & n = m \end{cases}$$

(19)

Note that we recover the original MWC model results, eqs 5–12, for the $n = 0$ wild-type ion channel. Similarly, the homooligomeric $n = m$ channel is also governed by the MWC model with $K_O \to K^\beta_0$ and $K_C \to K^\epsilon_0$. We also show the [EC_{50}] and $h$...
formulas for the \( n = \frac{m}{2} \) case to demonstrate the fundamentally different scaling behavior that this hetero-oligomeric channel exhibits with the MWC parameters.

Figure 8A shows that all of the CNGA2 mutants have small leakiness, which can be understood from their small \( \epsilon \) value and eq 16. In addition, the first term in the dynamic range eq 17 is approximately 1 because the open state affinities are always smaller than the closed state affinities by at least a factor of 10, which is then raised to the fourth power. Thus, all of the mutants have a large dynamic range as shown in Figure 8B.

Experimentally, single channel measurements confirmed that the wild-type \( n = 0 \) channel is nearly always closed in the absence of ligand; in the opposite limit of saturating cGMP, it was found that \( p_{\text{open}} \approx 0.99 \) for both the \( n = 0 \) and \( n = m \) ion channels (see Supporting Information section C.2).13

The \( [EC_{50}] \) and effective Hill coefficient \( h \) are shown in Figure 8C,D. In contrast to the nAChR case, where each additional mutation decreased \([EC_{50}]\), each CNGA2 mutation tends to increase \([EC_{50}]\), although not by a uniform amount. The effective Hill coefficient has a particularly complex behavior, first decreasing with each of the first three subunit mutations and then finally increasing back to the wild-type level for the fully mutated ion channel. To explain this decrease in \( h \) for the hetero-oligomeric channels, we first note that the wild-type \( n = 0 \) channel has a sharp response about its \([EC_{50}] \approx e^{-\beta\epsilon/mK_O} \) while the fully mutated \( n = m \) channel has a sharp response about \([EC_{50}] \approx e^{-\beta\epsilon/mK_B} \). Roughly speaking, the response of the hetero-oligomeric channels with \( 1 \leq n \leq 3 \) will occur throughout the full range between \( e^{-\beta\epsilon/mK_O} \) and \( e^{-\beta\epsilon/mK_B} \) which causes the dose-response curves to flatten out and leads to the smaller effective Hill coefficient. Such behavior could influence, for example, the response of the hetero-oligomeric nAChR ion channel if the two acetylcholine binding pockets diverged to have different ligand affinities.

Although we have focused on the particular mutants created by Wongsamitkul et al., it is straightforward to apply this framework to other types of mutations. For example, in Supporting Information section B.2 we consider the effect of modifying the \( K_O \) and \( K_C \) parameters of all four CNGA2 channels simultaneously. This question is relevant for physiological CNGA2 channels where a mutation in the gene would impact all of the subunits in the homo-oligomer, in contrast to the Wongsamitkul constructs where the four subunits were stitched together within a single gene. We find that when \( K_O \) and \( K_C \) vary for all subunits, the leakiness, dynamic range, and effective Hill coefficient remain nearly fixed for all parameter values, and that only the \([EC_{50}]\) scales linearly with \( K_C \) as per eq 11. In order to affect the other properties, either the gating energy \( \beta\epsilon \) or the number of subunits \( m \) would need to be changed.

**Extrapolating the Behavior of a Class of Mutants.** In this section, we explore how constant trends in both the nAChR and CNGA2 data presented above provide an opportunity to characterize the full class of mutants based on the dose-response curves from only a few of its members. Such trends may well be applicable to other ion channel systems,
enabling us to theoretically probe a larger space of mutants than what is available from current data.

First, we note that because the \( [EC_{50}] \) of the five nAChR mutants fell on a line in Figure 4C, we can predict the response of the entire class of mutants by only considering the dose-response curves of two of its members and extrapolating the behavior of the remaining mutants using linear regression. Experimentally, such a characterization arises because the L251S mutation acts nearly identically and independently across subunits to change the gating free energy of nAChR.12,17,23 This implies that mutating \( n \) subunits would yield an ion channel with gating energy

\[
\epsilon(\alpha) = \epsilon(0) + n\Delta \epsilon
\]

(20)

where \( \epsilon(0) \) is the wild-type gating energy and \( \Delta \epsilon \) is the change in free energy per mutation. This functional form is identical to the mismatch model for transcription factor binding, where each additional mutation-induced mismatch adds a constant energy cost to binding.27 Figure 9A demonstrates how fitting only the \( n = 0 \) and \( n = 4 \) dose-response curves (solid lines) accurately predicts the behavior of the \( n = 1, 2, \) and 3 mutants (dashed lines). In Supporting Information section D, we carry out such predictions using all possible input pairs.

We now turn to the CNGA2 ion channel where, once the \( K_{O}, K_{C}, K_{O*}, K_{C*}, \) and \( \epsilon \) parameters are known, the dose-response curve of any mutant can be predicted by varying \( n \) in eq 15. Figure 9B demonstrates that the wild-type ion channel (\( n = 0 \)) and the ion channel with only one mutated subunit (\( n = 1 \)) can accurately predict the dose-response curves of the remaining mutants. Supporting Information section D explores the resulting predictions using all possible input pairs.

**MWC Model Allows for Degenerate Parameter Sets.**

One critical aspect of extracting model parameters from data is that degenerate sets of parameters may yield identical outputs, which suggests that there are fundamental limits to how well a single data set can fix parameter values.25,28 This phenomenon, sometimes dubbed “sloppiness,” may even be present in models with very few parameters such as the MWC model considered here. Figure 10 demonstrates the relationship between the best-fit parameters within the nAChR and CNGA2 systems. For concreteness, we focus solely on the nAChR system.

After the value of \( K_{O} \) is fixed (to that shown on the \( x \)-axis of Figure 10A), the remaining parameters are allowed to freely vary in order to best fit the nAChR data. Although every value of \( K_{O} \in [10^{-11}–10^{-9}] \) M yields a nearly identical response curve in excellent agreement with the data (with a coefficient of determination \( R^2 > 0.96 \)), we stress that dissociation constants are rarely found in the range \( K_{O} \ll 10^{-10} \) M. In addition, a dissociation constant above the nanomolar range, \( K_{O} \gg 10^{-9} \) M, cannot fit the data well and is therefore invalidated by the
model. Thus, we may suspect the true parameter values will fall around the interval $K_O \in [10^{-10} - 10^{-7}] \ M$ for the nAChR system. $K_O$ could ultimately be fixed by measuring the leakiness $\epsilon$ (and thereby fixing $\beta_e$) for any of the ion channel mutants.

Two clear patterns emerge from Figure 10A: (1) The value of $K_C$ is approximately constant for all values of $K_O$, and (2) the five free energies all vary as $\beta_e^{(5)} = 2 \log(K_O) + n$ (constant). This suggests that $K_C$ and $e^{-\beta_e^{(2)} K_O}$ are the fundamental parameter combinations of the system. In Supporting Information section C.4, we discuss how these parameter combinations arise from our model.

We end by noting that the notion of sloppiness, while detrimental to fixing the physical parameter values of a system, nevertheless suggests that multiple evolutionary paths may lead to optimal ion channel phenotypes, providing another mechanism by which allostery promotes a protein’s capacity to adapt.39

■ DISCUSSION

There is a deep tension between the great diversity of biological systems and the search for unifying perspectives to help integrate the vast data that has built up around this diversity. Several years ago at the Institut Pasteur, a meeting was convened to celebrate the 50th anniversary of the allostery concept, which was pioneered in a number of wonderful papers in the 1960s and which since then has been applied to numerous biological settings.30–34 Indeed, that meeting brought together researchers working in physiology, neuroscience, gene regulation, cell motility, and beyond, making it very clear that allosteroy has great reach as a conceptual framework for thinking about many of the key macromolecules that drive diverse biological processes.

In this paper, we have built upon this significant previous work and explored how the Monod-Wyman-Changeux model can serve as a unifying biophysical framework for whole suites of ion channel mutants (see Figures 3 and 6). Specifically, we used two well-studied ligand-gated ion channels to explore the connection between mutations, the MWC parameters, and the full spectrum of dose-response curves which may be induced by those mutations. In addition, we have shown how earlier insights into the nature of “data collapse” in the context of bacterial chemotaxis and quorum sensing9,20 can be translated into the realm of ion channels. By introducing the Bohr parameter, we are able to capture the nonlinear combination of thermodynamic parameters which governs the system’s response.

For both the nAChR and CNGA2 ion channels, we showed that precise predictions of dose-response curves can be made for an entire class of mutants by only using data from two members of this class (Figure 9). In other words, the information contained in a single dose-response curve goes beyond merely providing data for that specific ion channel. Ultimately, because the total space of all possible mutants is too enormous for any significant fraction to be explored experimentally, we hope that a coupling of theory with experiment will provide a step toward mapping the relation between channel function (phenotype) and the vast space of protein mutations.

Moreover, we used the MWC model to determine analytic formulas for key properties such as the leakiness, dynamic range, $[EC_{50}]$, and the effective Hill coefficient, which together encapsulate much of the information in dose-response curves. These relationships tie into the extensive knowledge about phenotype–genotype maps,27,29,35 enabling us to quantify the trade-offs inherent in an ion channel’s response. For example, when modifying the ion channel gating free energy, the changes in the leakiness and $[EC_{50}]$ are always negatively correlated (Figure 4), whereas modifying the ligand binding domain will not affect the leakiness but may change the $[EC_{50}]$ (Figure 8 and Supporting Information section B.2). The ability to navigate between the genotype and phenotype of proteins is crucial in many bioengineering settings, where site-directed mutagenesis is routinely employed to find mutant proteins with specific characteristics (e.g., a low leakiness and large dynamic range).36–38

While general formulas for these phenotypic properties were elegantly derived in an earlier work,39 we have shown that such relations can be significantly simplified in the context of ion channels where $1 < e^{-\beta_e}\ll \left(\frac{K_C}{K_O}\right)^{\omega}$ (see eqs 9–12). This approximation is applicable for the range of parameters spanned by both the nAChR and CNGA2 systems, and we suspect it may hold for many other ion channels. These formulas provide a simple, intuitive framework to understand the effects of mutations. For example, they suggest the following: (1) Channel pore mutations that increase $\epsilon$ will exponentially increase the leakiness of the channel, although the constraint $1 < e^{-\beta_e}$ ensures that this leakiness will still be small. Ligand domain mutations are not expected to affect leakiness at all. (2) Channel pore mutations will exponentially decrease the $[EC_{50}]$ with increasing $\epsilon$, although this effect is diminished for ion channels with multiple subunits. For mutations in the ligand binding domain, the $[EC_{50}]$ will increase linearly with the dissociation constant $K_O$ between the ligand and the open ion channel (see Supporting Information section B.2). (3) Neither the dynamic range nor the effective Hill coefficient will be significantly perturbed by either type of mutation. (4) Transforming a homo-oligomeric ion channel into a hetero-oligomer can generate a significantly flatter response. For example, even though the CNGA2 channel composed of either all wild-type or all mutant subunits had a very sharp response, a channel composed of both subunits had a smaller effective Hill coefficient (see Figure 8D).

The framework presented here could be expanded in several exciting ways. First, it remains to be seen whether channel pore mutations and ligand binding domain mutations are completely independent, or whether there is some cross-talk between them. This question could be probed by creating a channel pore mutant (whose dose-response curves would fix its new $\tilde{\epsilon}$ values) and a ligand domain mutant (whose new $K^*_o$ and $K^*_C$ values would be characterized from its response curve), and then creating the ion channel with both mutations. If these two mutations are independent, the response of the double mutant can be predicted a priori using $\tilde{\epsilon}$, $K^*_o$, and $K^*_C$.

We also note that the MWC model discussed here does not consider several important aspects relating to the dynamics of ion channel responses. Of particular importance is the phenomenon of desensitization which significantly modifies an ion channel’s response in physiological settings.39–40 In addition, some ion channels have multiple open and closed conformations41–43 while other channels exhibit slow switching between the channel states.44 Exploring these additional complexities within generalizations of the MWC model would be of great interest.

Finally, we believe that the time is ripe to construct an explicit biophysical model of fitness to calculate the relative...
importance of mutation, selection, and drift in dictating the diversity of allosteric proteins such as the ion channels considered here. Such a model would follow in the conceptual footsteps laid in the context of fitness of transcription factors binding,8,27,35a,43 protein folding stability,8,8,44 and influenza evolution.49 This framework would enable us to make precise, quantitative statements about intriguing trends; for example, nearly all nAChR pore mutations appear to increase a channel’s leakiness, suggesting that minimizing leakiness may increase fitness.12 One could imagine that computing derivatives such as \( \frac{d \text{leakiness}}{d \epsilon} \), a quantity analogous to the magnetic susceptibility in physics, would be correlated with how likely an \( \epsilon \) mutation is to be fixed. The goal of such fitness functions is to map the complexity of the full evolutionary space (i.e., changes to a protein amino acid sequence) onto the MWC parameters and then determine how these parameters evolve in time. In this way, the complexity of sequence and structure would fall onto the very low dimensional space governed by \( \epsilon \), \( K_0 \), and \( K_C \).

### CONCLUSIONS

We have shown that the Monod-Wyman-Changeux (MWC) model of allostery can quantitatively account for the behavior of broad classes of mutant ligand-gated ion channels using the three biophysical parameters featured in that model: the free energy difference \( \epsilon \) between the closed and open conformations of the ion channel in the absence of ligand and the dissociation constants \( K_0 \) and \( K_C \) between the ligand and the open and closed ion channel, respectively. By examining nAChR and CNGA2 ion channel mutants, we demonstrated that pore mutations can be quantitatively understood as only affecting the \( \epsilon \) parameter while mutations in the ligand binding site only alter \( K_0 \) and \( K_C \) (Figures 3 and 6).

Building upon these insights, we derived simple analytic approximations for several important properties of an ion channel’s response, namely, the leakiness, dynamic range, \( [\text{EC}_{50}] \), and effective Hill coefficient (eqs 9–12 for nAChR, eqs 16–19 for CNGA2). These formulas are not limited in scope to nAChR and CNGA2, but rather reflect more general properties of ligand-gated ion channels. Utilizing these analytic relations, we quantified the spectrum of possible phenotypes achievable by mutating an ion channel pore or ligand binding domain (Figures 4 and 8).

In examining ion channels through the lens of the MWC model, we gained several intriguing insights into the dynamics of channel behavior. First, critical information about the leakiness and dynamic range of a channel’s response gets lost when reporting normalized current rather than absolute current (or when only considering the Hill limit). Second, when the physical parameters induced by a series of mutations exhibit a simple trend, the behavior of an entire class of mutants may be extrapolated from only a few of its members (Figure 9). Third, care must be taken when extracting the physical parameters that characterize the MWC model from data, as multiple degenerate parameter sets can often lead to nearly identical response curves (Figure 10). Finally, data from disparate mutants collapse onto a single master curve as a function of the Bohr parameter eq 14, which represents the combination of physical parameters that completely characterizes an ion channel’s response. In such cases, rather than thinking of each mutant as a story unto itself, we can instead categorize and understand mutants through simple one-parameter families.

### ASSOCIATED CONTENT

* Supporting Information

The following Supporting Information is available: The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcb.6b12672. This material is available free of charge via the Internet at http://pubs.acs.org/.

Details on aforementioned derivations and calculations (PDF)

Supplementary Mathematica notebook from which all of the plots can be reproduced, including the dose-response data for both the nAChR and CNGA2 systems (ZIP)

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**Notes**

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